A NOTE ON THE PERIODIC ORBITS OF A SELF EXCITED RIGID BODY

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ABSTRACT. The aim of the present paper is to study the periodic orbits of a perturbed self excited rigid body with a fixed point. For studying these periodic orbits we shall use averaging theory of first order.

1. INTRODUCTION AND STATEMENT OF THE MAIN RESULTS

The equations of the motion of a rigid body with moments of inertia $A, B, C$ $(A \leq B \leq C)$ under a external torque $\mathcal{M} = (M_1, M_2, M_3)$ are

$$
\begin{align*}
\frac{d\pi_1}{dt} &= \frac{B - C}{BC} \pi_2 \pi_3 + M_1, \\
\frac{d\pi_2}{dt} &= \frac{C - A}{AC} \pi_1 \pi_3 + M_2, \\
\frac{d\pi_3}{dt} &= \frac{A - B}{AB} \pi_1 \pi_2 + M_3,
\end{align*}
$$

where $\Pi = (\pi_1, \pi_2, \pi_3)$ is the total angular momentum of the rigid body and $M_i$ are smooth functions depending on the variables $\pi_1, \pi_2, \pi_3$. This rigid body is usually called at the literature self–excited rigid body, see for more information on this mechanical system Leimanis [9] or Deprit et alt. [2] In the following, we assume that we work with a rigid body quasi–axisymmetric under small torques, i.e. $B = A + \varepsilon$, and the functions $M_i = \varepsilon f_i(\pi_1, \pi_2, \pi_3)$ with $\varepsilon$ a small real parameter. Under these assumptions, very usual in many practical applications, the equations of motion (1) are reduced to

$$
\begin{align*}
\frac{d\pi_1}{dt} &= -\frac{C - A - \varepsilon}{(A + \varepsilon)C} \pi_2 \pi_3 + \varepsilon f_1(\pi_1, \pi_2, \pi_3), \\
\frac{d\pi_2}{dt} &= \frac{C - A}{AC} \pi_1 \pi_3 + \varepsilon f_2(\pi_1, \pi_2, \pi_3), \\
\frac{d\pi_3}{dt} &= -\frac{\varepsilon}{A(A + \varepsilon)} \pi_1 \pi_2 + \varepsilon f_3(\pi_1, \pi_2, \pi_3),
\end{align*}
$$

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The objective of this work is, by means of the Averaging Theory (see Theorem 4 of the Appendix) to provide a system of nonlinear equations whose simple zeros provide periodic solutions of (2). This idea has been used in previous works as an integrability criterium and for obtaining approximation solutions to real problems, see for instance [3, 4, 5, 6]. In order to present our results we need some preliminary definitions and notation.

We define the following functions in the open domain \( \Omega = \{ (R, Z) \in \mathbb{R}^2 \mid 0 < R < +\infty, 0 < Z < +\infty \} \) given by,

\[
F_1(R, Z) = \frac{1}{\pi Z} \int_0^{2\pi} \left( \cos \theta f_1 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) + \sin \theta f_2 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) \right) d\theta
\]

and \( F_2(R, Z) = \frac{2\alpha}{\pi Z} \int_0^{2\pi} f_3 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) d\theta, \)

with \( \alpha = \frac{C - A}{CA} \). A zero \( (R_0^*, Z_0^*) \) of the nonlinear system

\[
F_1(R, Z) = 0, \quad F_2(R, Z) = 0,
\]

such that

\[
\det \left( \frac{\partial (F_1, F_2)}{\partial (R, Z)} \right)_{(R, Z) = (R_0^*, Z_0^*)} \neq 0,
\]

is called a simple zero of system (4).

Our main results on the periodic solutions of the self-excited rigid body (2) are the following.

**Theorem 1.** Then for \( \epsilon \neq 0 \) sufficiently small and for every simple zero \( (R_0^*, Z_0^*) \neq (0, 0) \) of the nonlinear system (4), then the self-excited rigid body (2) has a periodic solution \( (\pi_1(t, \epsilon), \pi_2(t, \epsilon), \pi_3(t, \epsilon)) \) tending when \( \epsilon \to 0 \) to

\[
\left( \frac{R_0^*}{2\alpha}, 0, Z_0^* \right).
\]

Theorem 1 is proved in section 2. Its proof is based in the averaging theory for computing periodic solutions, see the appendix.

We provide two applications of Theorem 1 in the following two corollaries, which will be proved in section 3, and example of perturbed maps where Theorem 1 not provides information, see section?? for details.

**Corollary 2.** If \( f_1(\pi_1, \pi_2, \pi_3) = -\pi_3^2 \pi_1, \quad f_2(\pi_1, \pi_2, \pi_3) = \frac{2}{3} \left( \frac{\pi_3^3 \pi_2}{2} + \pi_3 \pi_2 \right) \)

and \( f_3(\pi_1, \pi_2, \pi_3) = \pi_1^2 - 2\pi_3 \pi_1 \)

then the differential equation (2) for \( \epsilon \neq 0 \) sufficiently small has two periodic solutions \( (\pi_1(t, \epsilon), \pi_2(t, \epsilon), \pi_3(t, \epsilon)) \) tending when \( \epsilon \to 0 \) to the following initial conditions

\[
(1, 0, 1) \quad \text{and} \quad (X_0^*, Y_0^*) = \left( \frac{1}{2}, 0, 2 \right),
\]
respectively.

**Corollary 3.** If $f_1(\pi_1, \pi_2, \pi_3) = -\alpha(\pi_1 + \pi_2)$, $f_2(\pi_1, \pi_2, \pi_3) = \alpha^3\pi_2^2$ and $f_3(\pi_1, \pi_2, \pi_3) = 1 - \alpha^2\pi_1^2\pi_3$ then the differential equation (2) for $\varepsilon \neq 0$ sufficiently small has one periodic solution $(\pi_1(t, \varepsilon), \pi_2(t, \varepsilon), \pi_3(t, \varepsilon))$ tending when $\varepsilon \to 0$ to the following initial condition $$\left(\frac{2\sqrt{3}}{3\alpha}, 0, \frac{3}{2}\right).$$

2. **Proof of the Theorem 1**

**Proof of the Theorem 1.** Introducing the change of variables

$$\pi_1 = \frac{1}{2\alpha}R \cos \theta, \quad \pi_2 = \frac{1}{2\alpha}R \sin \theta, \quad \pi_3 = Z,$$

in the system (2) by means of the following relations

$$\frac{dR}{dt} = \frac{\pi_1}{R} \frac{d\pi_1}{dt} + \frac{\pi_2}{R} \frac{d\pi_2}{dt}, \quad \frac{d\theta}{dt} = \frac{\pi_1}{R^2} \frac{d\pi_2}{dt} - \frac{\pi_2}{R^2} \frac{d\pi_1}{dt}, \quad \frac{dZ}{dt} = \frac{d\pi_3}{dt},$$

we obtain, to first order in the small parameter $\varepsilon$, the following system of differential equations

$$\begin{align*}
\frac{dR}{dt} &= \varepsilon F_1(R, \theta, Z) + O(\varepsilon^2), \\
\frac{d\theta}{dt} &= Z \frac{4\alpha}{4\alpha} + O(\varepsilon), \\
\frac{dZ}{dt} &= \varepsilon F_2(R, \theta, Z) + O(\varepsilon^2),
\end{align*}$$

with

$$F_1(R, \theta, Z) = \frac{\Delta}{4\alpha^2C^2},$$

$$F_2(R, \theta, Z) = f_3 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) - \frac{R^2(1 + \alpha C)^2 \sin 2\theta}{8\alpha^2C},$$

where $\Delta$ is equal to $2\alpha C^2 \cos \theta f_1 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) + 2\alpha C^2 \sin \theta$

$$f_2 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) + RZ(1 + \alpha C)^2 \cos \theta.$$

Using the equation $\frac{d\theta}{dt} = \frac{Z}{4\alpha} + O(\varepsilon)$ the system (5) is reduced to a planar system, $2\pi$–periodic in the variable $\theta$ in the following form

$$\begin{align*}
\frac{dR}{d\theta} &= \varepsilon \frac{4\alpha F_1(R, \theta, Z)}{Z} + O(\varepsilon^2), \\
\frac{dZ}{d\theta} &= \varepsilon \frac{4\alpha F_2(R, \theta, Z)}{Z} + O(\varepsilon^2),
\end{align*}$$

with

$$F_1(R, \theta, Z) = \frac{\Delta}{4\alpha^2C^2},$$

$$F_2(R, \theta, Z) = f_3 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) - \frac{R^2(1 + \alpha C)^2 \sin 2\theta}{8\alpha^2C},$$

where $\Delta$ is equal to $2\alpha C^2 \cos \theta f_1 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) + 2\alpha C^2 \sin \theta$

$$f_2 \left( \frac{R \cos \theta}{2\alpha}, \frac{R \sin \theta}{2\alpha}, Z \right) + RZ(1 + \alpha C)^2 \cos \theta.$$

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$$\begin{align*}
\frac{dR}{d\theta} &= \varepsilon \frac{4\alpha F_1(R, \theta, Z)}{Z} + O(\varepsilon^2), \\
\frac{dZ}{d\theta} &= \varepsilon \frac{4\alpha F_2(R, \theta, Z)}{Z} + O(\varepsilon^2),
\end{align*}$$
in the set \( \Omega_1 = \{(R, \theta, Z) \in \mathbb{R}^3 / 0 < R < +\infty, \theta \in S_1, 0 < Z < +\infty\} \). System (6) is in normal form for using, see the Appendix. Computing the averaged system we obtain

\[
\frac{dR}{d\theta} = \varepsilon F_1(R, Z) + O(\varepsilon^2),
\]

\[
\frac{dZ}{d\theta} = \varepsilon F_2(R, Z) + O(\varepsilon^2),
\]

in the open set \( \Omega = \{(R, Z) \in \mathbb{R}^2 / 0 < R < +\infty, 0 < Z < +\infty\} \) with \( F_1(R, Z) \) and \( F_2(R, Z) \) given by (3). Hence Theorem 1 is proved. \( \square \)

3. PROOF OF THE COROLLARY 2 AND 3

Proof of the corollary 2. Under the assumptions of Corollary 2 the nonlinear system (4) becomes

\[
F_1(R, Z) = \frac{R(2 - 3Z + Z^2)}{6\alpha},
\]

\[
F_2(R, Z) = \frac{Z}{2} \left( 8\alpha - \frac{2R^2Z^2}{\alpha} \right).
\]

This system has the two real solutions

\( (R_0^*, Z_0^*) = (2\alpha, 1) \) and \( (R_0^*, Z_0^*) = (\alpha, 2) \).

Moreover

\[
\det \left( \frac{\partial (F_1, F_2)}{\partial (R, Z)} \right)_{(R, Z) = (2\alpha, 1)} = \frac{4}{3}
\]

and

\[
\det \left( \frac{\partial (F_1, F_2)}{\partial (R, Z)} \right)_{(R, Z) = (\alpha, 2)} = \frac{8}{3}
\]

check that these solutions are simple. So, by Theorem 1 we only have two periodic solutions of (2). This completes the proof of the corollary. \( \square \)

Proof of the corollary 3. Under the assumptions of Corollary 3 the nonlinear system (4) becomes

\[
F_1(R, Z) = \frac{R(3R^2 - 16)}{Z},
\]

\[
F_2(R, Z) = \frac{\alpha(8 - R^2Z)}{2Z}.
\]

This system has the only solution

\( (R_0^*, Z_0^*) = (\frac{4\sqrt{3}}{3}, \frac{3}{2}) \).

Moreover

\[
\det \left( \frac{\partial (F_1, F_2)}{\partial (R, Z)} \right)_{(R, Z) = (R_0^*, Z_0^*)} = \frac{-32\alpha}{27} \neq 0
\]
check that these solutions are simple. So, by Theorem 1 we only have one periodic solutions of (2). This completes the proof of the corollary. □

4. Appendix

Now we shall present the basic results from averaging theory that we need for proving the results of this paper.

The next theorem provides a first order approximation for the periodic solutions of a periodic differential system, for the proof see [1] and [10].

Consider the differential equation
\[ \dot{x} = \varepsilon F_1(t, x) + \varepsilon^2 R(t, x, \varepsilon), \quad x(0) = x_0 \]
with \( x \in \mathbb{R}^n \), \( t \geq 0 \). Moreover we assume that both \( F_1(t, x) \) and \( F_2(t, x) \) are \( T \)-periodic in \( t \). Separately we consider in \( D \) the averaged differential equation
\[ \dot{y} = \varepsilon f_1(y), \quad y(0) = x_0, \]
where
\[ f_1(y) = \frac{1}{T} \int_0^T F_1(t, y) dt. \]

Under certain conditions, equilibrium solutions of the averaged equation turn out to correspond with \( T \)-periodic solutions of equation (8).

**Theorem 4.** Consider the two initial value problems (7) and (8). Suppose:

(i) \( F_1 \), its Jacobian \( \partial F_1 / \partial x \), its Hessian \( \partial^2 F_1 / \partial x^2 \), \( F_2 \) and its Jacobian \( \partial F_2 / \partial x \) are defined, continuous and bounded by an independent constant \( \varepsilon \) in \([0, \infty) \times D \) and \( \varepsilon \in (0, \varepsilon_0] \).

(ii) \( F_1 \) and \( F_2 \) are \( T \)-periodic in \( t \) (\( T \) independent of \( \varepsilon \)).

(iii) \( y(t) \) belongs to \( \Omega \) on the interval of time \([0, 1/\varepsilon]\).

Then the following statements hold.

(a) For \( t \in [1, \varepsilon] \) we have that \( x(t) - y(t) = O(\varepsilon) \), as \( \varepsilon \to 0 \).

(b) If \( p \) is a singular point of the averaged equation (8) and
\[ \det \left( \frac{\partial (f_1 + \varepsilon f_2)}{\partial y} \right) \bigg|_{y=p} \neq 0, \]
then there exists a \( T \)-periodic solution \( \varphi(t, \varepsilon) \) of equation (7) which is close to \( p \) such that \( \varphi(0, \varepsilon) \to p \) as \( \varepsilon \to 0 \).

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