FUNDAMENTALS OF PHYSICS IN ENGINEERING I

Degree in Sound and Image Engineering (in Telecommunications)
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SUMMARIES OF UNITS

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Fundamentals of Physics in Engineering I

Unit 1.- KINEMATICS

• Introduction
Mechanics is the branch of physics concerned with the relationships between the motion of physical systems and the forces that cause this motion. Mechanics is divided into three parts: Kinematics which studies motion without considering the forces that cause it; Dynamics which studies motion and the forces that cause it; and Statics which studies forces on bodies in static equilibrium.

• Position, velocity and acceleration
The first step in describing the motion of a particle is to establish a system of coordinates or reference system. The position vector, \( \mathbf{r} \), situates an object with respect to the origin of a reference system and, in general, is a function of time. In Cartesian coordinates:
\[
\mathbf{r}(t) = x(t) \mathbf{i} + y(t) \mathbf{j} + z(t) \mathbf{k}
\]
When a particle moves, the end of vector \( \mathbf{r} \) describes a curve called trajectory. If \( s \) is the space covered by the particle along its trajectory, \( s \) will be a function of time \( t \) and the function \( s = f(t) \) is known as a law of motion.
The displacement vector \( \Delta \mathbf{r} \) is the change in position vector between two points \( P_1 \) and \( P_2 \):
\[
\Delta \mathbf{r} = \mathbf{r}_2 - \mathbf{r}_1
\]
The average velocity \( \mathbf{v}_m \) of a particle is given by the displacement of a point during an interval of time \( \Delta t \), divided by the time interval:
\[
\mathbf{v}_m = \frac{\Delta \mathbf{r}}{\Delta t}
\]
The instantaneous velocity \( \mathbf{v} \) is the limiting value of the average velocity as the time interval \( \Delta t \) tends to zero. The following conditions are met:
\[
\mathbf{v} = \frac{d\mathbf{r}}{dt}
\]
The instantaneous velocity vector is tangent to the trajectory of the particle at each point of its trajectory.
The average acceleration \( \mathbf{a}_m \) of a material point is the change in velocity during an interval of time \( \Delta t \), divided by this time interval:
\[
\mathbf{a}_m = \frac{\Delta \mathbf{v}}{\Delta t}
\]
The instantaneous acceleration \( \mathbf{a} \) is the limit of average acceleration as the time interval tends to zero:
\[
\mathbf{a} = \frac{d\mathbf{v}}{dt} = \frac{d^2\mathbf{r}}{dt^2}
\]
• Intrinsic components of acceleration
The instantaneous acceleration \( \mathbf{a} \) may be decomposed into two vectors, one perpendicular to the trajectory called normal or centripetal acceleration, \( \mathbf{a}_N \), and the other, tangential to the trajectory, called tangential acceleration, \( \mathbf{a}_T \). These components are known as intrinsic components of acceleration:
\[
\mathbf{a} = \mathbf{a}_N + \mathbf{a}_T
\]
\( \mathbf{a}_T \) takes into account the change in modulus of velocity vector, \( \mathbf{v} = |\mathbf{v}| \), and \( \mathbf{a}_N \) takes into account the change in direction of velocity vector \( \mathbf{v} \):
\[
\mathbf{a}_N = \frac{d\mathbf{v}}{dt}, \quad \mathbf{a}_T = \frac{\mathbf{v}^2}{r}
\]
where \( r \) is the radius of curvature of the trajectory of the particle at each point of the trajectory. The following condition is met:
\[
\mathbf{a} = \sqrt{\mathbf{a}_N^2 + \mathbf{a}_T^2}
\]
• Linear motion
Linear motion is one in which the trajectory is a straight line and the space covered coincides with the modulus of the displacement vector. Moreover, the radius of curvature is infinite and there is no normal acceleration.
In uniform linear motion, the velocity is constant and so the acceleration is zero. If motion takes place along the \( X \) axis, the following conditions are met:
\[
a(t) = 0 \quad \mathbf{v}(t) = \mathbf{v} = \text{cte.} \quad x(t) = x_0 + \mathbf{v} t
\]
In the case of uniformly accelerated linear motion, the acceleration is constant and the following conditions are met:
\[
a(t) = a = \text{cte.} \quad \mathbf{v}(t) = \mathbf{v}_0 + at \quad x(t) = x_0 + \mathbf{v}_0 t + \frac{1}{2} a t^2
\]

• Circular motion
Circular motion is motion in a plane and the trajectory is a circumference of radius \( R \). The space covered, \( s \), may be written as a function of angle \( \theta \) as follows:
\[
\theta = \frac{s}{R}
\]
The angular velocity \( \omega \) is the variation in \( \theta \) with time \( t \):
\[
\omega = \frac{d\theta}{dt}
\]
The following condition is met:
\[
\omega = \frac{v}{R}
\]
The angular acceleration \( \alpha \) is the variation in angular velocity \( \omega \) with time \( t \):
\[
\alpha = \frac{d\omega}{dt} = \frac{d^2 \mathbf{r}}{d \theta^2}
\]
The following condition is met:
\[
\alpha = a \frac{\omega}{R}
\]
A vector \( \mathbf{\phi} \) may be assigned to the angular velocity \( \omega \) and another \( \mathbf{\alpha} \) to the angular acceleration \( \alpha \). In both cases the vectors are perpendicular to the circular trajectory of the particle and fulfill the following conditions:
\[
\mathbf{v} = \mathbf{\phi} \times \mathbf{r} \quad \mathbf{\alpha} = \mathbf{\phi} \times \mathbf{v} = \mathbf{\phi} \times (\mathbf{\omega} \times \mathbf{r})
\]
where \( \mathbf{r} \) is the vector that goes from the centre of the circumference to the position of the particle.
In uniform circular motion, the angular acceleration is zero and the angular velocity constant. This implies that there is no tangential acceleration (the modulus \( \mathbf{v} \) is also constant) and the normal acceleration is constant since \( \mathbf{v} \) and the radius \( R \) are constant. The following conditions are met:
\[
\alpha(t) = 0 \quad \omega(t) = \omega = \text{cte.} \quad \theta(t) = \theta_0 + \omega t
\]
In uniformly accelerated circular motion, the angular acceleration is constant. The tangential acceleration is constant but the normal acceleration is not. The following conditions are met:
\[
\alpha(t) = \alpha = \text{cte.} \quad \omega(t) = \omega_0 + \alpha t \quad \theta(t) = \theta_0 + \omega t + \frac{1}{2} \alpha t^2
\]

• Compositions of motion, Projectile motion
Another example of motion in a plane is the motion of a projectile launched at a constant velocity \( v_0 \) making an angle \( \alpha \) with the \( X \) axis and affected by the acceleration of gravity \( g \) along the \( Y \) axis. The trajectory is a parabola and the motion is composed of uniform linear motion along the \( X \) axis and uniformly decelerated linear motion along the \( Y \) axis. The time of flight, \( t \), maximum height, \( h \), and range, \( d \), are:
\[
t = \frac{2 v_0 \sin \alpha}{g}, \quad h = \frac{v_0^2 \sin^2 \alpha}{2g}, \quad d = \frac{v_0^2 \sin 2 \alpha}{g}
\]
The equation for the trajectory, \( y(x) \), is:
\[
y = \frac{g}{2v_0^2 \cos^2 \alpha} x^2 + x \tan \alpha
\]
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Unit 2.- DYNAMICS

• Introduction
Dynamics is the branch of mechanics concerned with motion and the forces that cause it. The motion of a body is a direct result of its interaction with other surrounding bodies and these interactions are conveniently described using the concept of force. The mass of a body is a measure of the object's resistance to changes in its velocity.

• Newton's laws
Newton's laws are fundamental laws of nature and form the basis of mechanics.

Newton's first law of motion (law of inertia): An object at rest remains at rest and an object in motion remains in motion at the same speed and in the same direction unless acted on by any kind of force.

Newton's second law of motion (fundamental equation of dynamics): the acceleration \( a \) of an object is directly dependent upon the net force \( F \) acting on the body:

\[
F = ma
\]

Newton's third law of motion (principle of action and reaction): If body A exerts a force \( F_{BA} \) (action) on body B, then body B exerts a force \( F_{AB} \) (reaction) on A of the same intensity but in the opposite direction. In other words, for every action there is an equal and opposite reaction:

\[
F_{AB} = -F_{BA}
\]

The forces of action and reaction act on different bodies. Newton's laws are only valid in an inertial reference frame, that is, a reference system in which an object at rest remains at rest if no net force acts upon it. Any reference system that moves at constant velocity with respect to an inertial system is also an inertial reference system. The surface of the earth is an approximate inertial reference system.

• Gravitational force, Weight
The Law of Gravitation was stated by Newton and makes it possible to determine the force of attraction \( F_{12} \) between two bodies of masses \( m_2 \) and \( m_2 \) separated by a distance \( r \):

\[
F_{12} = -\frac{G m_1 m_2}{r^2}
\]

where \( G = 6.67 \times 10^{-11} \text{ N m}^2 \text{ kg}^{-2} \) is the gravitational constant and \( u_r \) is a unitary vector in the direction of vector \( r \) that joins the two masses. The force of gravity is always one of attraction. Mass characterises two different properties of an object: its resistance to changes in its velocity (inertial mass) and its gravitational interaction with other objects (gravitational mass). Experiments show that the two are proportional and in the system of units chosen they are the same.

Supposing the earth has a radius \( R \) and mass \( M \), a body of mass \( m \) situated on the earth's surface will be attracted by a force \( F = GMm/R^2 \), and according to Newton's second law, such a mass will be subject to an acceleration \( g \):

\[
g = \frac{GM}{R^2}
\]

which is the acceleration due to gravity. The weight \( P \) of a body is the force exerted by the earth on the body:

\[
P = mg
\]

• Application of Newton’s laws to solve problems
The procedure for solving a problem in mechanics is as follows:
(i) Sketch the system and identify the object (or objects) to which Newton’s second law will be applied. Use vectors to represent the different forces.
(ii) Draw a diagram that includes the coordinate axes in order to decompose the vectors into their components. These diagrams should be drawn in such a way as to simplify the subsequent calculations. Normally this means including as many axes as possible along the directions of the forces, or placing an axis in the direction of acceleration, if this direction is known.
(iii) Using this diagram, write the components of Newton's second law as a function of known and unknown quantities and solve these equations for each unknown quantity. Finally, substitute the numerical values of the known quantities (including the units) in the equations and calculate each of the unknown quantities.

• Linear momentum and angular momentum
The linear momentum or amount of motion \( p \) of a particle of mass \( m \) moving at a velocity \( v \) is given by:

\[
p = mv
\]

Taking into account the relation \( a = dv/dt \), Newton's second law may be written as:

\[
F = \frac{dp}{dt}
\]

According to the law of conservation of linear momentum, in an isolated system, that is, one that is not subjected to external forces, the linear momentum is conserved. The mechanical impulse of a force \( J \) is defined as

\[
J = \int_{t_1}^{t_2} F dt
\]

The change in amount of motion of an object produced by a single impulsing force applied on the object is given by:

\[
\Delta p = p_2 - p_1 = J
\]

The angular momentum \( L \) of a particle of mass \( m \) with respect to a point \( O \) is given by:

\[
L = \mathbf{r} \times \mathbf{p}
\]

where \( \mathbf{r} \) is the vector which originates at the point \( O \) and ends at the particle, and \( \mathbf{p} = mv \) is the linear momentum of the particle. It may also be written as:

\[
L = m \mathbf{r} \times \mathbf{v}
\]

which shows that \( L \) is perpendicular to the velocity vector. Deriving the equation \( \mathbf{L} = \mathbf{r} \times \mathbf{p} \) with respect to time gives:

\[
\frac{dL}{dt} = \mathbf{r} \times F
\]

That is, the variation in angular momentum of a particle is equal to the momentum of the total force acting on the particle.

The law of conservation of angular momentum states that if the momentum of the total force acting on a particle is zero (\( \mathbf{r} \times \mathbf{F} = 0 \)), the angular momentum remains constant:

\[
\frac{dL}{dt} = 0 \Rightarrow L = \text{cte.}
\]

If the angular momentum remains constant, the trajectory of the particle is confined to a plane. In order for \( \mathbf{r} \times \mathbf{F} = 0 \) the following conditions must be met:
(i) \( \mathbf{F} = 0 \) (free particle)
(ii) \( \mathbf{F} \) and \( \mathbf{r} \) are two parallel vectors (central force).
• Introduction
Work and energy are two of the most important concepts in physics and also in everyday life. In physics, a force performs work when it acts on an object which moves a distance and a component of the force acts along the line of motion of the object. The concept of energy is closely linked to that of work. When a system performs work on another system, energy is transferred between the two systems. There are many forms of energy. Kinetic energy is the energy an object possesses due to its motion. Potential energy is energy which results from position or configuration, such as the distance between a body and the earth. Thermal energy is the kinetic energy resulting from the random movement of atoms and molecules within a system and is closely related to temperature. One of the fundamental laws of nature is the law of conservation of energy. If the energy of a system is conserved, its total energy remains constant, but some of it may be transformed from one form to another.

• Work and power
The work \( W \) performed by a force \( F \) that acts on a body moving along a trajectory is defined by the integral:

\[
W = \int \mathbf{F} \cdot d\mathbf{r}
\]

In the simple case of a constant force and distance \( \Delta r \) along a straight line, the work is given by the scalar product:

\[
W = F \cdot \Delta r
\]

For a force that varies along one dimension (for example, along the \( X \) axis):

\[
W = \int_{x_1}^{x_2} F_x(x) \cdot dx
\]

The unit of work in the SI is the Joule (J). The power \( P \) is the speed at which a force performs work:

\[
P = \frac{dW}{dt}
\]

The power of a force \( F \) performing work on an object moving at a velocity \( v \) is given by:

\[
P = F \cdot v
\]

In the SI, power is measured in watts (W).

• Kinetic energy. Kinetic energy theorem
The kinetic energy \( E_c \) of a body of mass \( m \) moving at a velocity \( v \) is:

\[
E_c = \frac{1}{2} mv^2
\]

Kinetic energy is the energy resulting from motion. The kinetic energy theorem states that the work done by a force acting on a body is equal to the change in kinetic energy of the body:

\[
W = \frac{1}{2} mv_{\text{final}}^2 - \frac{1}{2} mv_{\text{initial}}^2 = E_{c,\text{final}} - E_{c,\text{initial}}
\]

that is:

\[
W = \Delta E_c
\]

• Conservative forces and potential energy
A force is conservative if the work done in moving a particle around a closed trajectory is zero. In addition the work done by a conservative force is independent of the path taken and depends solely on the initial and final points. The work done by the weight of a body near the earth's surface is:

\[
W = -mg(y_2 - y_1)
\]

and it is independent of the trajectory connecting the initial and final points. Such a force is conservative. The potential energy \( E_p \) depends solely on position. Two examples of potential energy are gravitational potential energy:

\[
E_p = mgy
\]

and elastic potential energy due to compression or stretching of a spring:

\[
E_p = \frac{1}{2} k x^2
\]

In the case of a conservative force, the work \( W \) and potential energy \( E_p \) are related by the equation:

\[
W = -\Delta E_p
\]

and the force \( F \) and potential energy \( E_p \) by the equation:

\[
F = -\nabla E_p
\]

which in the one-dimensional case may be written as:

\[
F_x = -\frac{dE_p}{dx}
\]

The motion of an object may be represented by a potential energy graph. The points of equilibrium may be identified on this graph.

• Conservation of mechanical energy
The sum of the kinetic and potential energy of a system is called mechanical energy \( E \):

\[
E = E_c + E_p
\]

If there are no external forces acting on the system and all the internal forces are conservative, the total mechanical energy of the system remains constant:

\[
E = E_{c,1} + E_{p,1} = E_{c,2} + E_{p,2}
\]

The total energy of the system \( E_{\text{sist}} \) is the sum of the different types of energy. One way of transferring the energy (absorbed or ceded) of a system is to exchange work with the surroundings. If this is the only source of transferred energy, the law of conservation of energy is written as:

\[
W_{\text{ext}} = \Delta E_{\text{sist}}
\]

\( W_{\text{ext}} \) is the work done on the system by external forces and \( \Delta E_{\text{sist}} \) is the variation in total energy of the system. This is the work-energy theorem.
Unit 4.- HEAT AND TEMPERATURE

- Introduction
  Temperature is often said to be a measure of the degree of hotness or coldness of a body, but this definition is not valid from the point of view of physics. When a body heats up or cools down, some of its physical properties change. For example, most solids and liquids dilate when heated, when an electrical conductor is heated its resistance varies, etc. A physical property that varies with temperature is called a thermometric property and a change in this property indicates that the temperature of the object has changed.

- Thermal equilibrium and zeroth law of thermodynamics
  Two systems in contact are in thermal equilibrium when their properties no longer change with time. For two systems to be in contact must be separated by a diathermal wall that facilitates their thermal interaction. An adiabatic wall does not allow such interaction: each system is isolated from the other and each may remain in its equilibrium state. The zeroth law of thermodynamics states that if two systems are in thermal equilibrium with a third system, they are also in thermal equilibrium with each other. The concept of temperature is related to the state of thermal equilibrium of two systems since they will be in thermal equilibrium if they have the same temperature.

- Thermometers and ideal gas temperature scale
  In order to establish a temperature scale a thermometric property is used. Gas thermometers are characterized by the fact that they all register the same temperature provided that the density of the gas used in the thermometer is very low. The temperature of an ideal gas is defined using a limit with real gases diluted in a constant volume gas thermometer. The temperature scale is adjusted taking a temperature of 273.16 K as the triple point of water. In this state, the melting point, boiling point and sublimation point coincide and it occurs at a vapour pressure of 610 Pa and temperature of 0.01°C. The ideal gas temperature is defined as:

  \[ T = \frac{p}{\rho_3} \]

  where \( p \) is the pressure of the gas in the thermometer when it is in thermal equilibrium with the system whose temperature is to be measured, and \( \rho_3 \) is the pressure when the thermometer is in a ice-water-steam bath at its triple point.

- Ideal gas law
  The ideal gas equation is:

  \[ pV = nRT \]

  where \( p \) is the pressure, \( V \) the volume, \( n \) the number of moles, \( T \) the absolute temperature and \( R \) the ideal gas constant, whose numerical value depends on the units of the other physical magnitudes in the equation. In SI units, \( p \) is in Pa, \( V \) in m³, \( n \) in moles and \( T \) in K:

  \[ R = 8.31 \text{ J mol}^{-1} \text{ K}^{-1} \]

  If \( p \) is expressed in atmospheres (atm) and \( V \) in litres (l)

  \[ R = 0.082 \text{ l atm mol}^{-1} \text{ K}^{-1} \]

- Thermal dilation
  Most substances expand or dilate when their temperature rises and contract when it falls.
  (i) Linear dilatation linear: The change \( \Delta L \) in length \( L_0 \) of an object due to a change in temperature \( \Delta T \) is:

  \[ \Delta L = \alpha L_0 \Delta T \]

  where \( \alpha \) is the coefficient of linear dilatation of the substance.

  (ii) Surface dilatation: The change \( \Delta S \) in the surface \( S_0 \) of an object due to a change in temperature \( \Delta T \) is given by:

  \[ \Delta S = \gamma S_0 \Delta T \]

  where \( \gamma \) is the coefficient of surface dilatation of the substance. For an isotropic substance \( \gamma = 2\alpha \).

  (iii) Cubic dilatation: The change \( \Delta V \) in volume \( V_0 \) of an object due to a change in temperature \( \Delta T \) is given by:

  \[ \Delta V = \beta V_0 \Delta T \]

  where \( \beta \) is the coefficient of cubic dilatation or volume expansion. In the case of an isotropic substance, \( \beta = 3\alpha \). Water exhibits an anomalous thermal expansion between 0° and 4°C, since it contracts when the temperature rises.

- Quantity of heat: heat capacity and specific heat
  When a cold spoon is put into a cup of hot coffee, the spoon heats up and the coffee cools down in order to approach thermal equilibrium. These changes in temperature are basically caused by a transfer of energy from one substance to the other. The transfer of energy that occurs solely due to a difference in temperature is called heat flow or heat transfer, and the energy transferred is called heat. The symbol \( Q \) is used to denote the quantity of heat.

  A calorie is the quantity of heat necessary to raise the temperature of 1 g of water from 14.5°C to 15.5°C. The relation between a joule and a calorie is: 1 cal = 4.186 J. The quantity of heat \( dQ \) supplied to a system of mass \( m \) and the change in temperature produced \( dT \) are related by the specific heat \( c \) which depends on the material:

  \[ dQ = mc \, dT \]

  The quantity of heat \( Q \) needed to raise the temperature of a mass \( m \) of a material from \( T_1 \) to \( T_2 \) is approximately proportional to the change in temperature \( \Delta T = T_2 - T_1 \) and mass \( m \) of the material

  \[ Q = mc \Delta T \]

  If the number of moles \( n \) is known we can refer to the molar heat capacity \( C \):

  \[ dQ = nC \, dT \quad C = \frac{1}{n} \frac{dQ}{dT} \]

  where \( n = m/M \), and \( M \) is the molecular mass. We can write:

  \[ Q = nC \Delta T \]

  In the case of solids, we normally work with specific heat and molar heat capacity at constant pressure \( c_p \) and \( C_p \), whereas with gases, constant volume \( cv \) and \( C_v \).

- Calorimetry, phase changes and latent heat
  The heat necessary to melt a solid substance is given by:

  \[ Q_f = mL_f \]

  where \( L_f \) is the latent heat of fusion. For water at atmospheric pressure \( L_f = 333.5 \text{ kJ/kg} = 80 \text{ cal/g} \). The heat necessary to vaporize a liquid is given by:

  \[ Q_v = mL_v \]
where \( L_v \) is the latent heat of vaporization. For water at atmospheric pressure \( L_f = 2257 \text{ kJ/kg} = 540 \text{ cal/g} \). Under the same conditions, each phase change occurs at a certain temperature, and while a phase change is taking place the temperature of the system remains constant.

• Propagation of heat by conduction

The heat \( Q \) is the energy transferred between a system and its surroundings due solely to a difference in temperature between the system and a part of its surroundings. The heat flow continues until the temperatures are the same. When heat is propagated due to conduction, it is transmitted between two systems by direct contact. If the means separating the systems, whose temperatures are \( T_1 \) and \( T_2 \), has a length \( L \) and section \( S \), in the stationary state (\( T \) no longer changes with time) the heat that passes through a transversal section per unit of time (thermal current, \( H = Q/t \)) is given by:

\[
H = k S \frac{T_2 - T_1}{L}
\]

where \( k \) is the thermal conductivity of the medium. The thermal resistance of the medium, \( R \), is given by:

\[
R = \frac{L}{kS}
\]

and \( H = \Delta T/R \). For a compound wall in the stationary state, its equivalent thermal resistance is the sum of the thermal resistances of the wall components if they have the same surface area.

The thermal current \( H \) in the case of non-stationary conditions and diverse geometries is given by:

\[
H = -k S \frac{dT}{dx}
\]

which is known as Fourier's law. \( dT/dx \) is the temperature gradient and \( H \) is the instantaneous heat current through an element of area \( S \). The negative sign indicates that the heat flows from high to low temperatures.

• Propagation of heat by convection and radiation

When heat is propagated by convection, heat is transferred from one place to another by movement of heated material giving rise to macroscopic convection currents, which may appear in liquids in a gravitational field whose density varies with temperature (natural convection). Convection may also be forced using ventilators.

When heat is transferred by radiation, the power \( P \) radiated by a surface is given by Stefan-Boltzmann's law:

\[
P = \varepsilon \sigma S T^4
\]

where \( \varepsilon \) is the emissivity and \( \sigma = 5.67 \times 10^{-8} \text{ Wm}^{-2}\text{K}^{-4} \) is the Stefan-Boltzmann constant. All objects emit energy from their surfaces when they are hot and thermal radiation is a type of electromagnetic radiation.

In all types of heat propagation, if the temperature difference between the body and its surroundings is small, the body's rate of cooling is approximately proportional to the temperature difference (Newton's law of cooling).
**Fundamentals of Physics in Engineering I**

**Unit 6. - ELECTRIC FIELD**

- **Introduction**
  The ancient Greeks were the first to observe the phenomenon of electrical attraction. They found that when amber was rubbed, it attracted small objects such as straw or feathers. In fact, the word “electric” comes from the Greek word for amber, elektron.

  Electrical nature of matter. Electric charge
  Electric charge is a fundamental property of matter and two types of charge exist: positive and negative. Two bodies with the same type of charge repel each other, whereas if they are of opposite charge they attract each other.

  Quantification and conservation of electric charge
  The electric charge is always a multiple of the elementary charge or quantum of electric charge, whose value is given by:
  \[ e = 1.602177 \times 10^{-19} \text{ C} \]
  which is the absolute value of the charge on an electron.

  In all processes that take place in nature, the total or net charge of an isolated system remains constant.

- **Coulomb’s law: Electric force between point charges**
  Coulomb’s law expresses the electric force \( F \) exerted by a point charge \( q \) on another point charge \( q' \):
  \[ F = K_e \frac{qq'}{r^2} u_r \]
  where \( r \) is the vector starting at \( q \) and ending at \( q' \) and \( u_r = r/r \).
  \( K_e \) is the constant:
  \[ K_e = \frac{1}{4\pi\epsilon_0} = 9 \times 10^9 \text{ N C}^2 \text{ m}^2 \]
  This force is inversely proportional to the square of the distance, and is a force of attraction between charges of opposite signs and of repulsion between charges of the same sign.

- **Electric field**
  Electric field of a point charge
  If an electric charge experiences a force due to the presence of other charges in the region, an electric field is created. The electric field \( E \) generated by a charge distribution is the force \( F \) exerted by the distribution on a test particle divided by the charge \( q \) of the test particle:
  \[ E = \frac{F}{q} \quad F = qE \]
  For a point charge:
  \[ E = K_e \frac{q}{r^2} u_r \]
  where \( u_r \) is the unit vector that points along the line from the charge \( q \) to the point where the field \( E \) is evaluated.

  Principle of superposition
  The electric field created by a system of point charges is equal to the sum of the electric fields produced by each individual charge.

  Electric field lines
  The spatial characteristics of an electric field may be illustrated by drawing lines of force or electric field lines. The tangent of a line of force at each point indicates the direction of the field \( E \) at that point.

  The electric field lines start at the positive charges and end at the negative charges. A uniform field has the same intensity and direction at all points in space and is represented by rectilinear field lines that are parallel and equidistant.

  Electric field created by a continuous charge distribution
  For a continuous charge distribution (in volume, surface or line) the electric field is calculated as follows:
  \[ E = K_e \int \frac{dq}{r^2} u_r \]
  This may be applied, for example, to determine the field created by a ring or a disk at points on its axes or the field created by a rectilinear segment at points on its perpendicular bisector.

  Movement of charge in an electric field
  If an electric force is the only force acting on a particle of mass \( m \) and charge \( q \), according to Newton's second law the acceleration is given by \( a = qE/m \). When a particle moves in a uniform electric field, its motion is described by the kinematics of motion with constant acceleration. It is of particular interest to study the motion of a charge that enters a region where there is a uniform electric field perpendicular to the velocity vector of the charge.

- **Potential electric energy and electric potential**
  The electric force between two point charges is directed along a line joining the two charges and is inversely proportional to the square of the distance between them, as is the force of gravity between two masses. Like the force of gravity, the electric force is conservative. Thus, there is a potential energy function associated with the electric force. If a test point charge is placed in an electric field, its potential energy is proportional to its charge. The electric potential energy per unit of charge is a function of the position of the charge in space and is called electric potential.

  The electric force is conservative. The potential energy of a test particle in the field created by various fixed particles \( q_i \) is given by:
  \[ E_p = K_e \sum q_i \frac{q_j}{r_{ij}} \]
  (taking the origin of potential energy at infinity).

  The electric potential of a charge \( q \) is defined as:
  \[ V = \frac{E_p}{q} \quad E_p = qV \]
  The SI unit of potential is called one volt (V):
  \[ 1 \text{ V} = 1 \text{ J/C} \]

  Electric potential of a point charge
  For a point charge (with the origin of potential at infinity):
  \[ V = K_e \frac{q}{r} \]
  For a system of point charges (with the origin of potential at infinity):
  \[ V = K_e \sum \frac{q_i}{r_{ij}} \]

  Electric potential of a continuous charge distribution
  For a continuous charge distribution:
where $V = k_e \int \frac{dq}{r}$.

Applying this equation, it is possible to calculate the potential of, for example, a ring charged along its axis.

**Potential difference**

The difference in potential $\Delta V$ between two points 1 and 2 is related to the work $W$ done by the electric field:

$$W = -\Delta E_p = E_{p1} - E_{p2} = q(V_1 - V_2) = -q\Delta V$$

Given:

$$\Delta V = V_2 - V_1 = -\int_1^2 E \cdot dl$$

The potential difference $V_1 - V_2$ is the negative work per unit of charge done by the electric field on a positive test charge when the latter moves from point 1 to point 2. The potential difference $\Delta V$ is also the positive work per unit of charge that must be done against the electric field in order to move the charge from 1 to 2.

**Relation between electric potential and electric field**

The relation between electric field $E$ and electric potential $V$ is given by:

$$dV = -E \cdot dl$$

The electric field lines point in the direction in which the potential decreases. If the expression for $E$ is known, the potential $V$ at a point $P$ may be obtained by means of the line integral of $E$:

$$V = -\int L E \cdot dl$$

If $V$ is known, the field $E$ may be found by means of the gradient of $V$:

$$E = -\nabla V = -\frac{dV}{dx}$$

If the electric field is constant along a direction (for example, along the $X$ axis):

$$E_x = -\frac{dV}{dx}$$

If the potential depends solely on the module of $r$ (that is, $r$):

$$E = \frac{dV}{dr}$$

**Equipotential surfaces**

Surfaces that have the same electric potential at their points, that is $V = constant$, are known as equipotential surfaces. The field lines are perpendicular to equipotential surfaces. For a point charge, equipotential surfaces are spherical surfaces concentric with the charge.

**Electric flux. Gauss’s law**

The electric field flux over a surface $S$ is defined as the surface integral of the electric field vector over the whole surface:

$$\Phi_E = \int_S E \cdot dS$$

When the flux across a closed surface is calculated, this surface is known as a Gaussian surface. Field lines may be used to visualize the flux across a surface. The total flux may be positive, negative or zero. When it is positive, the flux is outgoing and when negative, incoming.

Gauss’s law states that the total electric flux through any closed surface is equal to the total (net) electric charge inside the surface divided by $\varepsilon_0$:

$$\Phi_E = \frac{q}{\varepsilon_0} = \int_E dS$$

**Calculation of electric field using Gauss’s law**

Gauss’s law may be used to find the electric field produced by highly symmetrical charge distributions such as infinite lines, planes or spheres. The crucial step in this process is to select the Gaussian surface.

**Properties of conductors**

Experimentally it is shown that certain substances have the property of allowing the movement of electric charges through them, whereas others prevent this flow. The former are called conductors and the latter insulators or dielectrics. Metals and their alloys, aqueous solutions of acids, bases and salts, etc are conductors; however, only metallic conductors will be studied. These are made up of positive ions occupying fixed positions in a crystalline lattice and electrons that have become detached from the metal ions and circulate freely and randomly within the conductor, giving rise to a kind of “electron gas” which causes the phenomenon of electric conduction.

**General properties of conductors in electrostatic equilibrium**

(i) The electric field inside a conductor in electrostatic equilibrium is zero

(ii) The net electric charge of a conductor in electrostatic equilibrium is found on its surface.

(iii) The electric field on the surface of a conductor in electrostatic equilibrium is perpendicular to the surface.

(iv) The surface of a conductor in electrostatic equilibrium is an equipotential surface.

**Electric field in the proximities of a conductor in electrostatic equilibrium**

The electric field at points near the surface of a conductor is perpendicular to the surface and is given by:

$$E = \frac{\sigma}{\varepsilon_0}$$

which is known as Coulomb's theorem. This is the field created by the total charge of the conductor and may be considered to be the sum of the field created by a small disk of area $dS$ and that created by the rest of the conductor. The field $E_1$ created by a small disk at nearby points is:

$$E_1 = \frac{\sigma}{2\varepsilon_0}$$

then the field $E_2$ created by the rest of the surface charges on the conductor is:

$$E_2 = \frac{\sigma}{2\varepsilon_0}$$

therefore, $E = E_1 + E_2$.

**Conductors in an electric field**

When a conductor is placed in an electric field, the field inside the conductor must be cancelled in order for the conductor to be in electrostatic equilibrium. This results in the charges in the conductor being reordered to create an electric field inside the conductor to compensate the applied electric field.

**Dielectric breakdown and the point effect**

Many non-conducting materials are ionized in very high electric fields and become conductors. This phenomenon is known as dielectric breakdown and the dielectric limit of an insulator is the maximum electric-field magnitude, $E_{max}$, that can exist in this material without producing dielectric breakdown. This value of the field is also known as the dielectric strength. When dielectric breakdown takes place the molecules of the material are ionized and the material starts to conduct. In a gaseous material such as air $E_{max} \approx 3 \times 10^6$ V/m and this effect is accompanied by the emission of light due to the recombination of electrons with the ionized molecules, a phenomenon known as arc discharge.

When a conductor is non-spherical in shape, its surface is equipotential but the surface charge density and the electric field just outside the conductor varies from one point to another. The electric field is stronger at the points near the regions of smaller radius of curvature of the conductor (point effect). If the conductor has very small radius of curvature points, the dielectric surrounding it may break down at relatively low potentials. Lightning conductors extract the charge from nearby clouds before the potential of the cloud reaches a destructively high value.
• Capacitance
A conductor with charge \( Q \) and potential \( V \), has a capacitance:

\[ C = \frac{Q}{V} \]

• Capacitors
A capacitor is an electric device used in circuits to store charge and electric energy. It is made up of two conductor plates with a potential difference \( \Delta V \) between them. The capacitance of a capacitor is given by:

\[ C = \frac{Q}{\Delta V} \]

In the SI capacity is measured in farads (1 F = 1 C/V). The capacity depends on the geometric design of the capacitor and the nature of the dielectric between its plates or within its housing.

For a parallel-plate capacitor in vacuum (a capacitor with two planoparellel conducting plates, each with area \( S \), separated by a distance \( d \)), the capacitance is:

\[ C = \varepsilon_0 \frac{S}{d} \]

Capacitors in series and parallel
The equivalent capacitance of various capacitors connected together is the capacitance of a single capacitor which, when used on its own, produces the same external effect. The equivalent capacitance of various capacitors in series is:

\[ \frac{1}{C_{eq}} = \frac{1}{C_1} + \frac{1}{C_2} + ... + \frac{1}{C_i} \]

When various capacitors are in parallel the equivalent capacitance of the system is calculated using the equation:

\[ C_{eq} = C_1 + C_2 + ... + C_i \]

• Electrostatic energy
The potential electric energy \( U \) stored in a charged capacitor is calculated as the work required to charge it:

\[ U = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} CV^2 = \frac{1}{2} QV \]

When this energy is associated with an electric field \( E \), the electric energy density \( u_E \) in the space occupied by the field (in a vacuum) is:

\[ u_E = \frac{1}{2} \varepsilon_0 E^2 \]

In a material medium, \( \varepsilon_0 \) may be substituted by \( \varepsilon \). The total electric-field energy \( U \) in a volume \( V \) may be calculated using the integral:

\[ U = \int_V u_E dV \]

• Dielectrics

Electrostatic properties of dielectrics
When a dielectric is placed within the housing of a capacitor in which there was a vacuum between the plates, the capacitance increases so that:

\[ C = \varepsilon_r C_0 \]

whereas the potential difference and electric field decrease:

\[ V = \frac{V_0}{\varepsilon_r} \quad E = \frac{E_0}{\varepsilon_r} \]

where \( \varepsilon_r = \varepsilon / \varepsilon_0 \).

There are polar and non-polar dielectrics. The molecules of the former have no electric dipolar moment, whereas those of the latter have a permanent electric dipolar moment.

When a non-polar dielectric is placed in an electric field, its atoms or molecules become electric dipoles oriented in the direction of the electric field. If the dielectric is polar, its permanent dipolar moments are oriented in parallel to the external field.

When the electric dipoles of a substance line up spontaneously (ferroelectric substances) or due to the action of an external electric field, the substance is said to be polarized.

A polarized dielectric has charges on its surface and, unless the polarization is uniform, in its volume as well. These polarization charges, however, are frozen in the sense that they are bound to atoms or molecules and are not free to move in the dielectric. In a conductor, the charges are capable of moving freely and are called free charges.
• Introduction
A conductor is a material which allows charged particles to move freely through it and these particles carry the charge on the conductor. For example, a metal may be considered to be a structure of positive ions occupying fixed positions on a lattice with the free electrons moving between them. The charge of all the free electrons is equal and opposite to the charge of all the ions, thus giving rise to a neutral material. The free electrons can move freely within the lattice of ions and carry the charge through the metal.

• Current and charge motion. Current density
An electric current is a flow of charged particles. The intensity $I$ of the electric current characterises the charge flowing through a circuit element:

$$ I = \frac{dQ}{dt} $$

The direction of the current is the same as the direction of the drift velocity $v$ of the positive charge carriers. In the SI, the intensity $I$ is expressed in amperes (A). The current density $J$ is the charge flow at a point inside a conducting medium:

$$ J = n q v $$

where $n$ is the number of charge carriers per unit of volume (concentration of particles) and $v$ is the drift velocity of the charge carriers. Successive collisions with the ions in the lattice cause the drift velocity of the charge carriers to be a constant average. If $j$ is uniform, its module $j$ is given by:

$$ j = \frac{I}{S} $$

and in the general case:

$$ I = \int J \cdot dS $$

Conventionally, the direction of the electric current is taken to be that of the positive charge carriers. When the charge carriers are negative, the direction of the current is contrary to the direction of movement of these negative charge carriers.

• Ohm’s law. Conductivity, resistivity and resistance
For a conductor at a constant temperature, Ohm’s law states that the quotient between the potential difference $V$ between two points on a conductor and the electric current $I$ circulating through the conductor is a constant known as electrical resistance, $R$:

$$ \frac{V}{I} = R \quad \Rightarrow \quad V = RI $$

In this case the conductor is said to be ohmic. In the SI, the resistance $R$ is expressed in ohms ($\Omega$).

For a conductor of length $l$ and section $S$, Ohm’s law may be written as:

$$ J = \frac{l}{RS} E = \frac{1}{\rho} E = \alpha E $$

where $\alpha = 1/\rho$ is the conductivity of the material and $\rho$ its resistivity. $\rho$ is expressed in $\Omega^{-1} m^{-1}$ and:

$$ R = \frac{\rho l}{S} $$

The resistivity of metals increases with temperature. Ohm’s law may be written in general as:

$$ J = \alpha E $$

and the drift velocity of the electrons may be written as:

$$ v = \frac{\alpha}{en} E $$

• Resistances in series and in parallel
The equivalent resistance of resistances connected in series is given by:

$$ R_{eq} = \sum R_i $$

and of resistances connected in parallel:

$$ \frac{1}{R_{eq}} = \sum \frac{1}{R_i} $$

• Energy aspects of an electric current. Joule’s law
A supply of energy is necessary to maintain an electric current since the charges must be accelerated by an electric field. The energy per unit of time, or power, required to maintain a current is given by $P = IV$. In the SI, power $P$ is expressed in watts (W). For conductors that obey Ohm’s law $V = RI$ and so:

$$ P = IV = IR^2 = \frac{V^2}{R} \quad (\text{Joule’s law}) $$
Types of solids

There are various types of bonds between atoms; however, in order to systematize their study two main types of bonds are considered: (a) ionic (formed between a metal and a non-metal) and (b) covalent (formed between non-metals). In practice, no substance forms totally ionic or totally covalent bonds and these two types are extreme cases.

From the macroscopic point of view a solid is a rigid or elastic substance, that is, a substance that has an elastic behaviour under the influence of either hydrostatic forces or strain and tension.

Considering their structure, solids may be divided into two types: amorphous and crystalline. Solids may also be classified based on the predominant type of bond between the atoms or ions making up the crystals. In this classification, five types may be considered: covalent, ionic, those with hydrogen bonds, molecular and metallic.

Energy bands. Conductors, insulators and semiconductors

The concept of energy band is very useful for understanding various properties of solids such as electric conductivity.

When atoms join together to form a solid, their external energy levels overlap giving rise to bands. At absolute zero, insulators and semi-conductors have a full valence band separated by a prohibited band of energy from an empty conduction band.

In conductors, the occupied energy bands - valence and conduction bands (allowed bands) - are separated by prohibited bands that can not be occupied.

Only partially filled bands can give rise to electric currents in a solid when an electric field is applied. In the theory of bands, the difference between conductors and insulators thus resides in whether or not there are any partially occupied bands.

Silicon and germanium have the same external structure as carbon; thus they would be expected to have an energy band diagram similar to that of diamond. However, since the energy gap $E_G$ between the valence and conduction bands is small, when the temperature is raised some electrons in the valence band may acquire sufficient energy to jump to the conduction band. This enables an electric current to be created when an electric field is applied, as in the case of a conductor. Therefore, such substances are called semiconductors.

The distinction between an insulator and a semiconductor resides solely in the value of $E_G$. At ambient temperature, $E_G$ is 1.12 eV for Si and 0.72 eV for Ge. There are two types of current carriers in a semiconductor, electrons and holes, and the total current is the sum of the currents due to each type of carrier.

Intrinsic and extrinsic semiconductors

In pure semiconductors (intrinsic) conduction takes place due to the electrons that are present simply because the material is purely crystalline (for example, Ge or Si) and not due to the presence of impurities. The only way in which conduction may take place is to provide the electrons with an energy equal to or greater than $E_G$ - thus enabling them to jump to the conduction band. This may be done by light or thermal excitation.

When minuscule quantities, of the order of one part per million, of the right impurities (different atoms) are added to a pure semiconductor, the latter may exhibit electrical conductivity over a greater range of temperatures (extrinsic semiconductor). This is known as doping. Thus, the electrical properties of a semiconductor may change drastically when small concentrations of donor impurities ($n$-type semiconductor) or acceptor impurities ($p$-type semiconductor) are added.

Semiconductor equation and electrical neutrality

Semiconductor equation

Let $n$ and $p$ are the concentrations of free electrons and holes in a semiconductor, respectively. In an intrinsic semiconductor, and due to thermal excitation, an electron moves out of the valence band and leaves behind a hole, the number of free electrons is equal to the number of holes

$$ n = p = n_i = p_i $$

where $n_i$ and $p_i$ are the intrinsic concentrations of carriers. In general, for any semiconductor, the processes of electron-hole pair generation and recombination are continuous, and an equilibrium is reached at each temperature in which the product $np$ is constant. Thus

$$ np = n_i^2 = p_i^2 = \text{constant} \quad \text{(for fixed } T \text{ and } E_G \text{)}$$

This equation, known as the semiconductor equation or mass action law, is essential when studying semiconductors and semiconducting devices and is valid for both intrinsic and extrinsic semiconductors in thermal equilibrium. When there is no equilibrium, that is, when $n$ and $p$ are governed by external conditions, the above equation is not valid.

The condition of electrical neutrality

If a semiconductor is doped, the mass action law is not sufficient to determine the concentration of carriers. In addition, it is necessary to include an expression that
links concentration and density of donor or acceptor impurities. This expression is given by the electrical neutrality condition:

\[
[\text{positive charges}] = [\text{negative charges}]
\]

If \(N_a\) denotes the concentration of acceptor impurities and \(N_d\) that of donor impurities (number of atoms per unit volume), the electrical neutrality condition is written as:

\[
p + N_d = n + N_a
\]

which is easy to understand considering the following points: (a) The semiconductor is electrically neutral if no external field acts upon it. (b) When donor and acceptor impurities are added, additional electrons and holes are present. (c) All the donor and acceptor impurities are ionized.

**Transport phenomena in semiconductors**

Charge transport phenomena may occur in a semiconductor due to either application of electric fields or the existence of carrier concentration gradients, that is, when the concentration varies depending on the position in the semiconductor material. There are numerous types of transport phenomena but only the following will be considered:

(a) Electric conduction due to charge transport caused by applying a uniform electric field.
(b) Electric conduction due to diffusion of charge caused by the existence of a carrier gradient.
(c) The Hall effect, whereby application of a magnetic field gives rise to an electric field (that is, a potential difference).

Since there are two types of charge carriers in a semiconductor (electrons and holes), the expression for displacement (or drift) current has two terms, one for electrons and one for holes:

\[
\vec{J} = e(n\mu_n + p\mu_p)\vec{E}
\]

and so the conductivity of semiconductors is given by:

\[
\sigma = e(n\mu_n + p\mu_p)
\]

**Semiconductor devices**

When a type \(p\) semiconductor is joined to a type \(n\) semiconductor a junction diode (p-n junction) is formed which is the basis for manufacturing semiconductor devices such as diodes or transistors. These devices play a fundamental role in modern electronics.

A diode contains a p-n junction, whereas a bipolar junction transistor contains two p-n junctions that may be p-n-p or n-p-n.